

# Hyperon-Quark Mixed Phase in Compact Stars

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- Non uniform structure and EOS of high density matter.
- Structure and mass of compact stars.

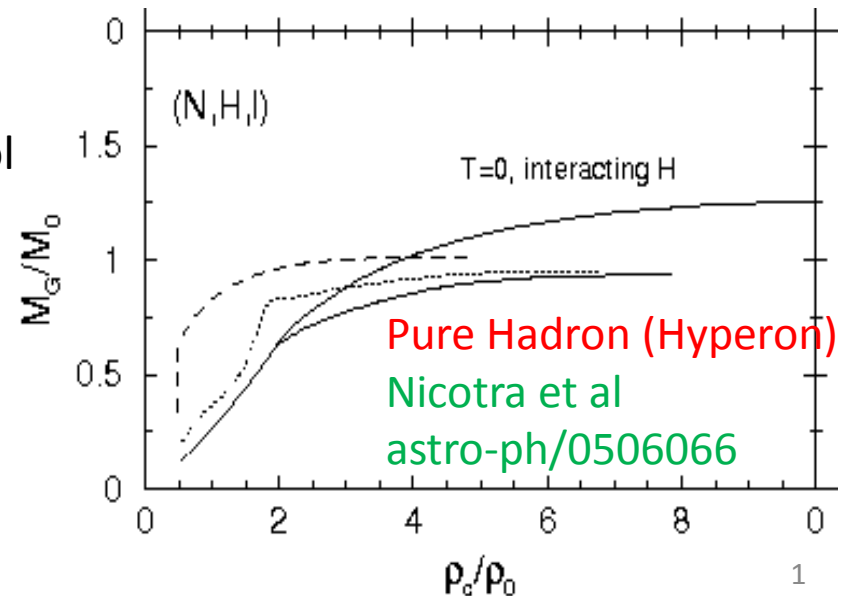
At 2-3  $\rho_0$ , hyperons are expected to appear.

→ Softening of EOS

→ Maximum mass of neutron star becomes

less than 1.4 solar mass.

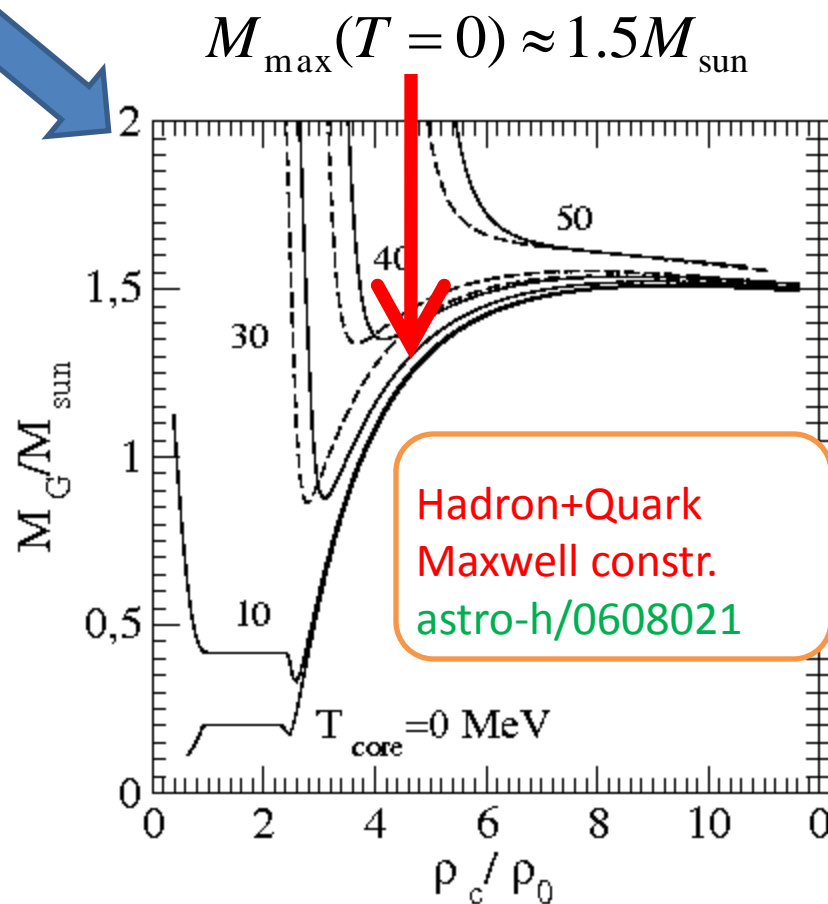
→ Contradicts the obs  $>1.5 M_{\text{sol}}$



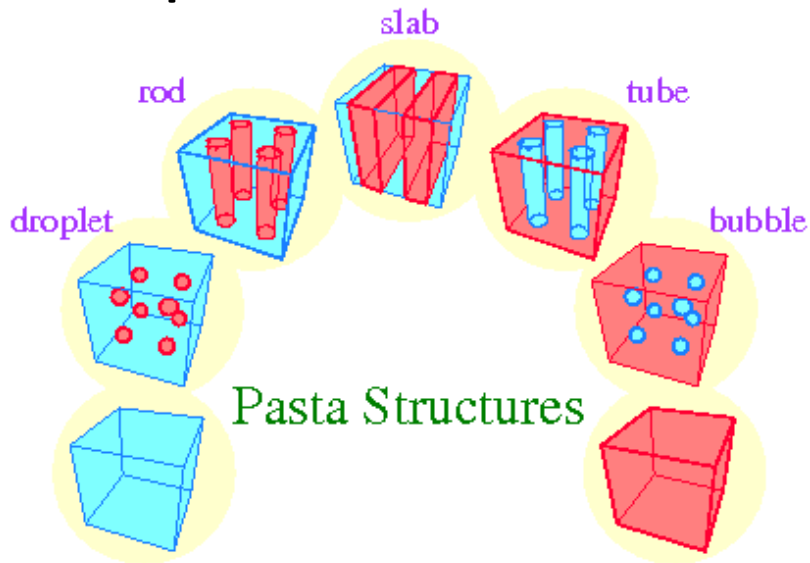
Phase transition to quark matter may solve this problem. [Catania group, PLB562,153 etc].

**But** the existence of mixed phase may soften the EOS again!

Bulk Gibbs calculation yields wide range of mixed phase and large softening [Glendenning, PRD46,1274].



In the mixed phase with charged particles, non-uniform “Pasta” structures are expected.



Depending on the density, geometrical structure of mixed phase changes from droplet, rod, slab, tube and to bubble configuration.

Quite different from a bulk picture of mixed phase. We have to take into account the effect of the structure when we calculate EOS.

# Hadron phase

## Brueckner Hartree Fock model

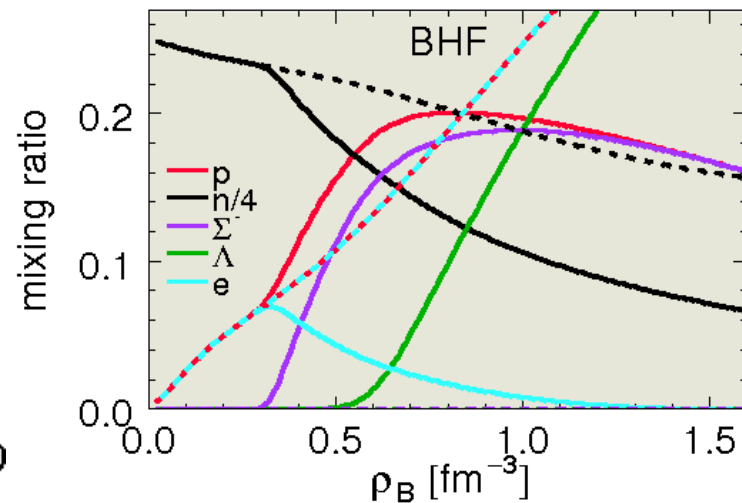
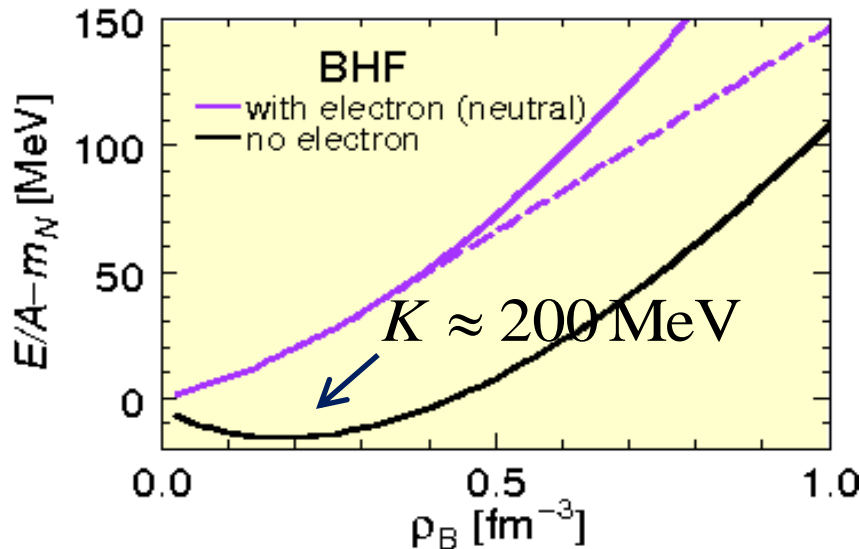
$$\varepsilon = \frac{3}{5} \frac{k_F^2}{2m} \rho + \frac{1}{2} \sum_{k, k' \leq k_F} \langle kk' | \mathcal{A} G[\rho; e(k) + e(k')] | kk' \rangle \quad (\mathcal{A} : \text{antisymmetrizer})$$

$$G[\rho; w] \equiv v + \sum_{k_a, k_b} v \frac{|k_a k_b \rangle Q \langle k_a k_b|}{w - e(k_a) - e(k_b)} G[\rho; w]$$

$$e(k) = e(k; \rho) \equiv \frac{k^2}{2m} + \text{Re} \sum_{k' \leq k_F} \langle kk' | \mathcal{A} G[\rho; e(k) + e(k')] | kk' \rangle$$

$v$  : AV18 + UIX + NSC89

YN scatt data,  $\Lambda$ -nuclear levels, and nuclear saturation are fitted.



# Quark phase

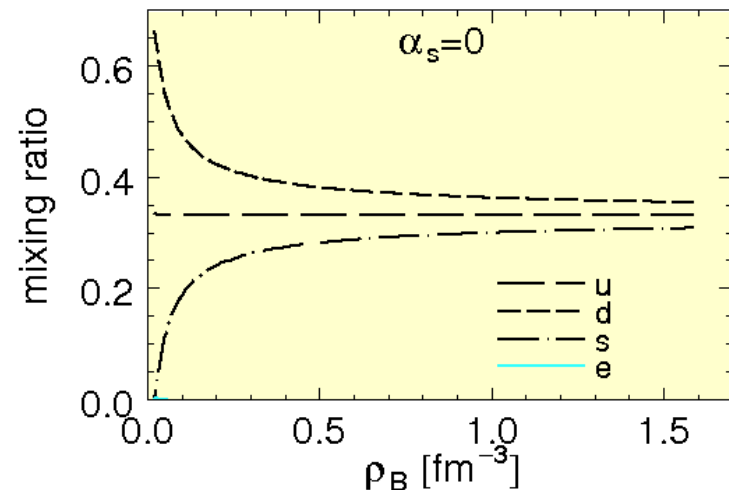
MIT bag model  $\varepsilon = \sum_f [\varepsilon_f^{\text{kin}} + \varepsilon_f^{\text{Fock}}] + \mathcal{B}$  ( $\mathcal{B}$ : bag constant)

$$\varepsilon_f^{\text{kin}} = \frac{3}{8\pi^2} m_f^4 \left[ x_f \eta_f (2x_f^2 + 1) - \log(x_f + \eta_f) \right]$$

$$\varepsilon_f^{\text{Fock}} = \frac{\alpha_s}{8\pi^3} m_f^4 \left[ x_f^4 - \frac{3}{2} \left[ x_f \eta_f - \log(x_f + \eta_f) \right]^2 \right]$$

$$x_f \equiv p_F^{(f)} / m_f, \quad \eta_f \equiv \sqrt{1 + x_f^2}$$

Electron fraction is very small in quark matter.

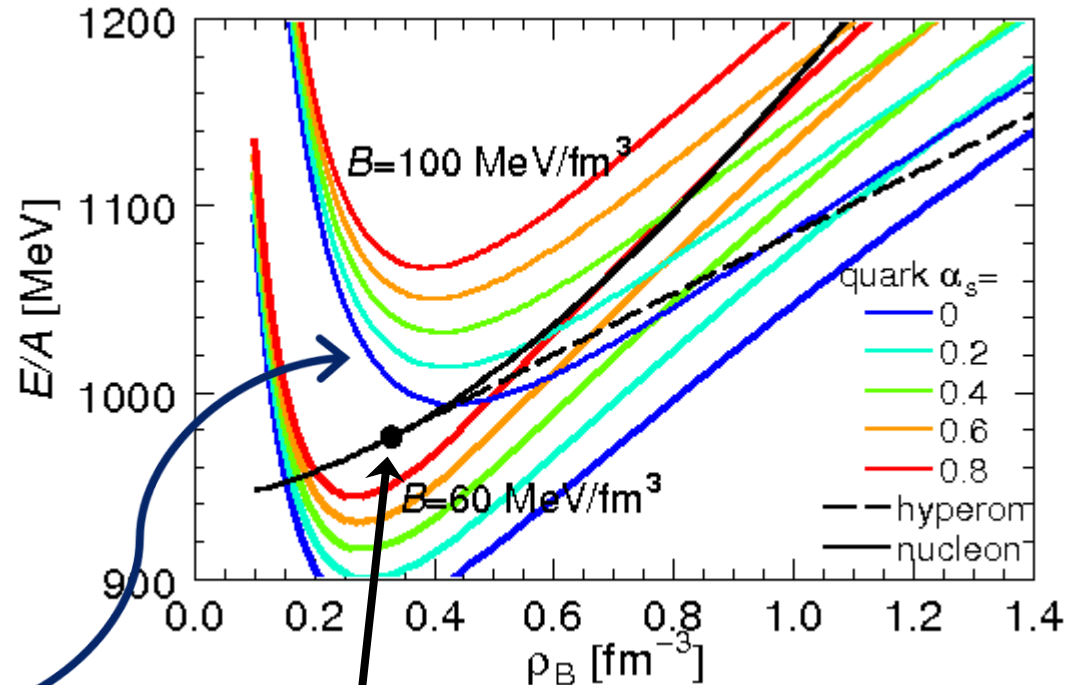


# Hadron EOS vs quark EOS

Depending on  $B$  and  $\alpha_s$ , hadron and quark EOSs cross at various densities.

We choose  $\alpha_s = 0$  and  $B = 100 \text{ MeV/fm}^3$ .

Quark threshold is above the hyperon threshold in uniform matter.



Hyperon threshold  
 $0.34 \text{ fm}^{-3}$

# Coupled equations

to get density profile, energy, pressure, etc of the system

$$\mu_u + \mu_e = \mu_d = \mu_s, \quad \mu_n = \mu_u + 2\mu_d, \quad \mu_p + \mu_e = \mu_n = \mu_\Lambda = \mu_\Sigma - \mu_e$$

$$\mu_i = \frac{\partial \varepsilon(\mathbf{r})}{\partial \rho_i(\mathbf{r})} \quad (i = u, d, s, p, n, \Lambda, \Sigma^-, e)$$

$$\varepsilon(\mathbf{r}) \equiv \varepsilon_B(\mathbf{r}) + \varepsilon_e(\mathbf{r}) + (\nabla V_C(\mathbf{r}))^2 / 8\pi e^2$$

$$\varepsilon_B(\mathbf{r}) = \begin{cases} \varepsilon_H(\mathbf{r}) & \text{(hadron phase)} \\ \varepsilon_Q(\mathbf{r}) & \text{(quark phase)} \end{cases}$$

$$\varepsilon_e(\mathbf{r}) = (3\pi^2 \rho_e(\mathbf{r}))^{4/3} / 4\pi^2$$

$$E/A = \frac{1}{\rho_B V} \left[ \int_V d^3 r \varepsilon(\mathbf{r}) + \sigma \mathcal{S} \right] \quad \left( \begin{array}{l} \rho_B = \text{average baryon density} \\ \mathcal{S} = \text{Q-H boundary area} \\ V = \text{cell volume} \end{array} \right)$$

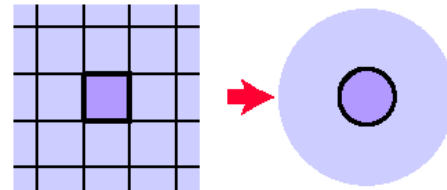
$$\int_V d^3 r \left[ \rho_p(\mathbf{r}) - \rho_\Sigma(\mathbf{r}) + \frac{2}{3} \rho_u(\mathbf{r}) - \frac{1}{3} \rho_d(\mathbf{r}) - \frac{1}{3} \rho_s(\mathbf{r}) - \rho_e(\mathbf{r}) \right] = 0 \quad (\text{total charge})$$

$$\frac{1}{V} \int_V d^3 r \left[ \rho_p(\mathbf{r}) + \rho_n(\mathbf{r}) + \rho_\Lambda(\mathbf{r}) + \rho_\Sigma(\mathbf{r}) + \frac{1}{3} \rho_u(\mathbf{r}) + \frac{1}{3} \rho_d(\mathbf{r}) + \frac{1}{3} \rho_s(\mathbf{r}) \right] = \rho_B \quad (\text{given})$$

# Numerical calculation

- Assume regular structures: Divide space into equivalent and charge-neutral cells with a geometrical symmetry (3D: sphere, 2D : cylinder, 1D: plate).

→ Wigner Seitz cell approximation



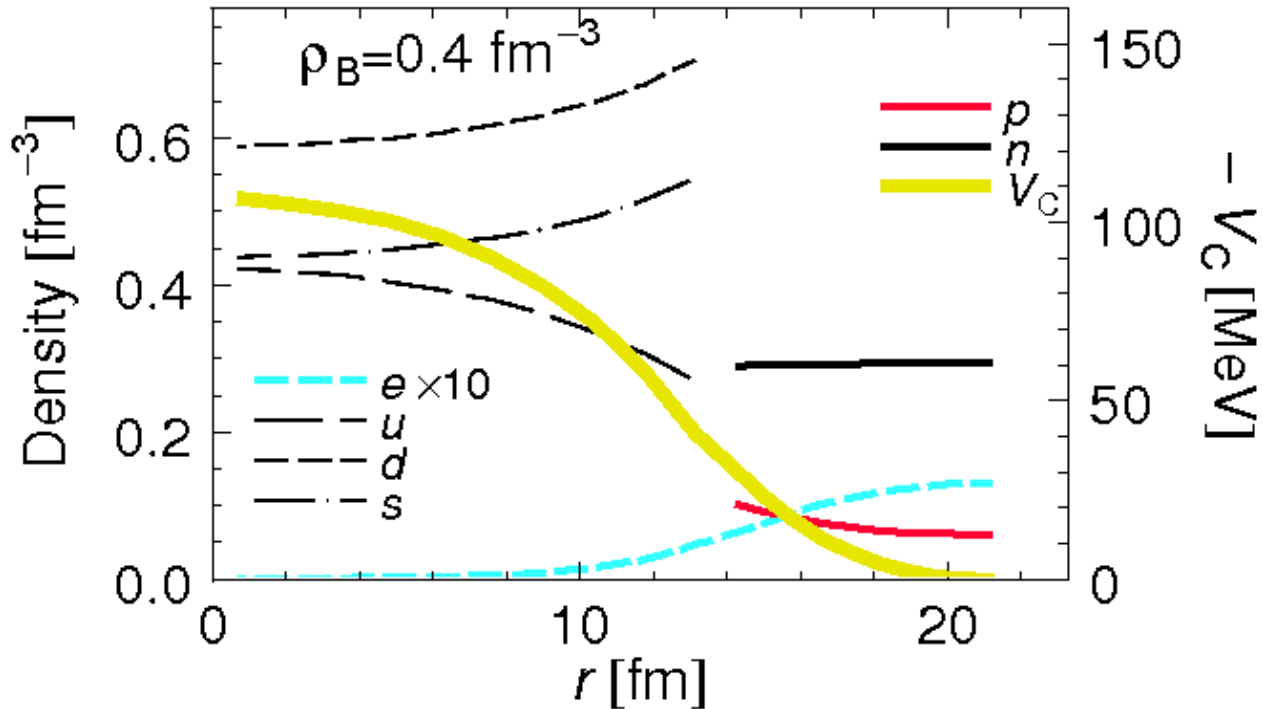
- Put a Q-phase, a H-phase and a phase boundary in the cell.



- Divide the cell into grid points and solve the field equations with **a given Baryon density**. Cell radius and boundary position are optimized.
- Compare 7 cases (Uniform H, 3D Q-droplet, 2D Q-rod, 1D slab, 2D tube, 3D bubble, Uniform Q) and choose the energy minimum solution.



# Density profile in a cell

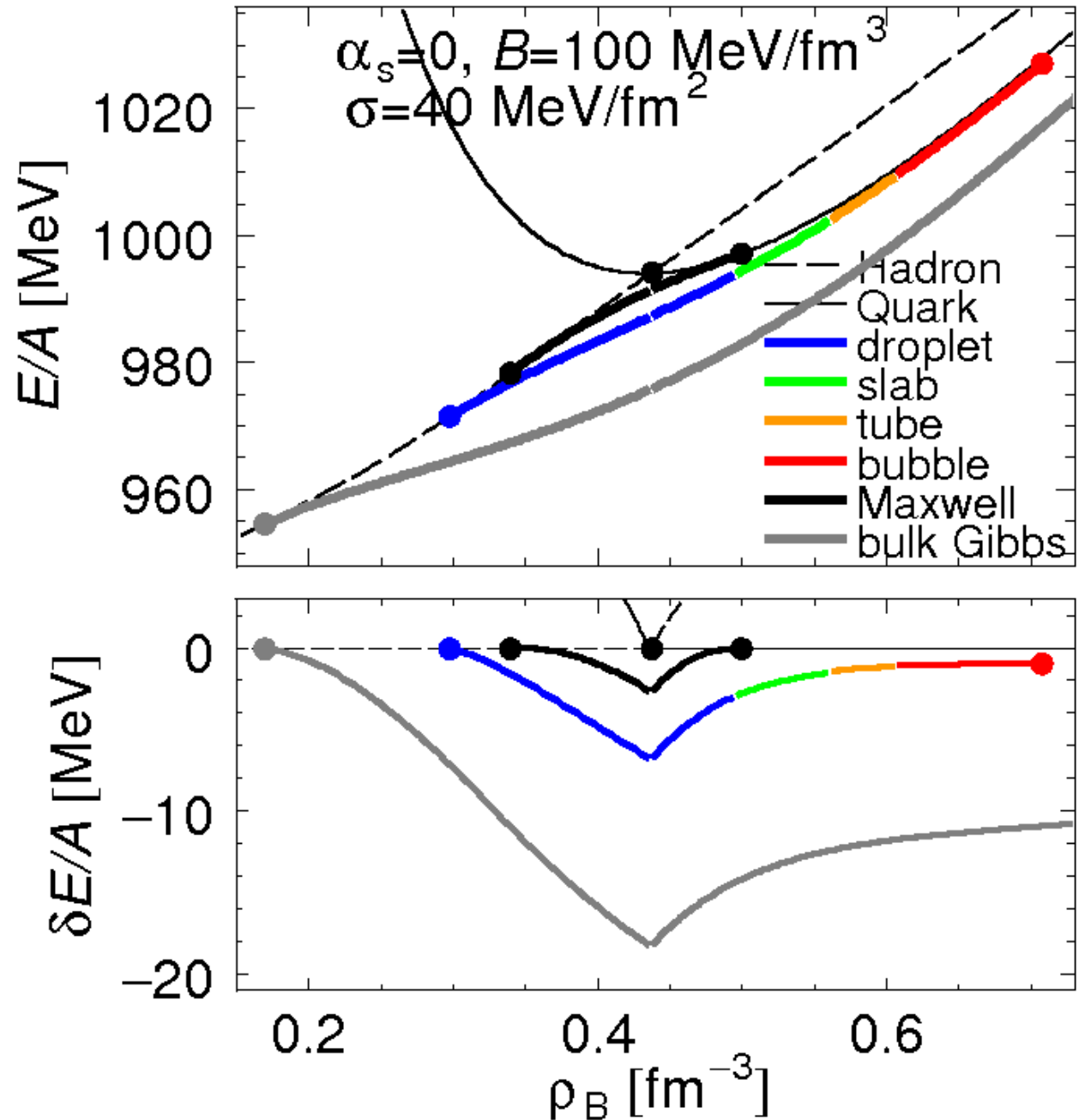


Quark phase is negatively charged.

→  $u$  quarks are attracted and  $ds$  quarks repelled.  
 Same thing happens to  $p$  in the hadron phase.

# EOS of matter

Full calculation is close to the **Maxwell construction** (local charge neutral). Far from the **bulk Gibbs** calculation (neglects the surface and Coulomb).



# Particle fraction

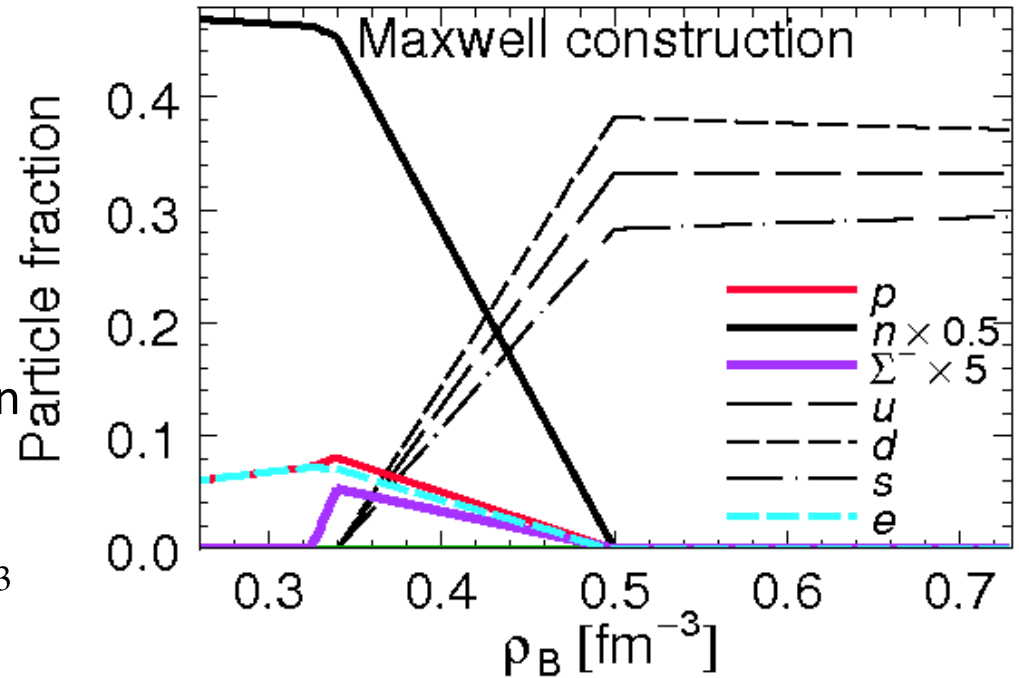
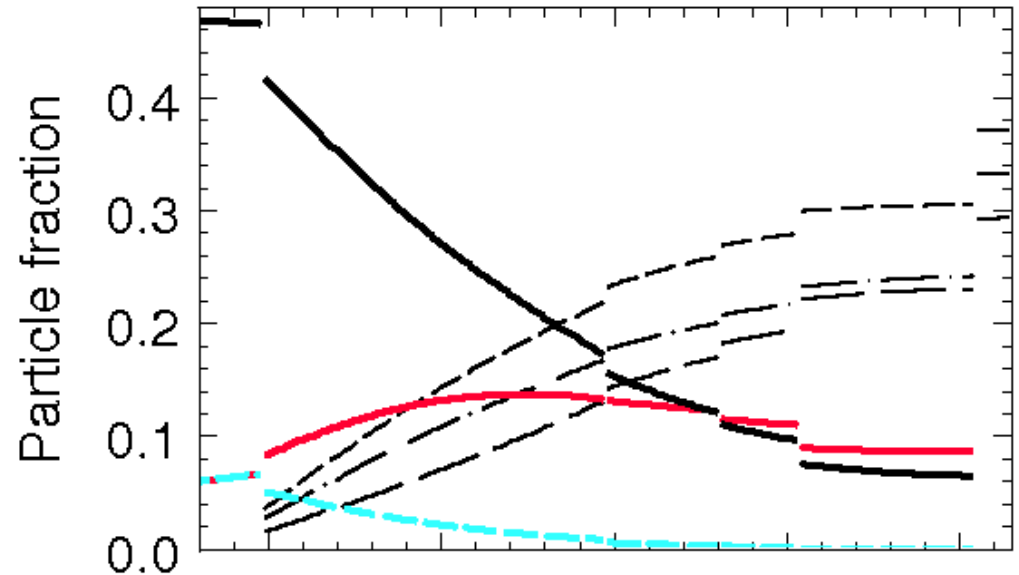
Though the EOS of full calc is close to the Maxwell constrctn, the particle fraction is much different.

Hyperon does not appear.

Lack of the charge-neutrality condition in each phase suppresses hyperons.

Neutral matter  $\rightarrow \rho_{th}=0.34 \text{ fm}^{-3}$

Non neutral matter  $\rightarrow \rho_{th}=1.15 \text{ fm}^{-3}$



# Structure of compact stars

Solve TOV eq.

$$\frac{dP}{dr} = -\rho \frac{Gm}{r^2} \left( 1 + \frac{4\pi r^3 P}{m} \right) \left( 1 + \frac{P}{\rho} \right) \left( 1 - \frac{2Gm}{r} \right)^{-1}$$

$$P = P(\rho)$$

Pressure (input of TOV eq.)

$$\rho = \rho(r)$$

Density at  $r$

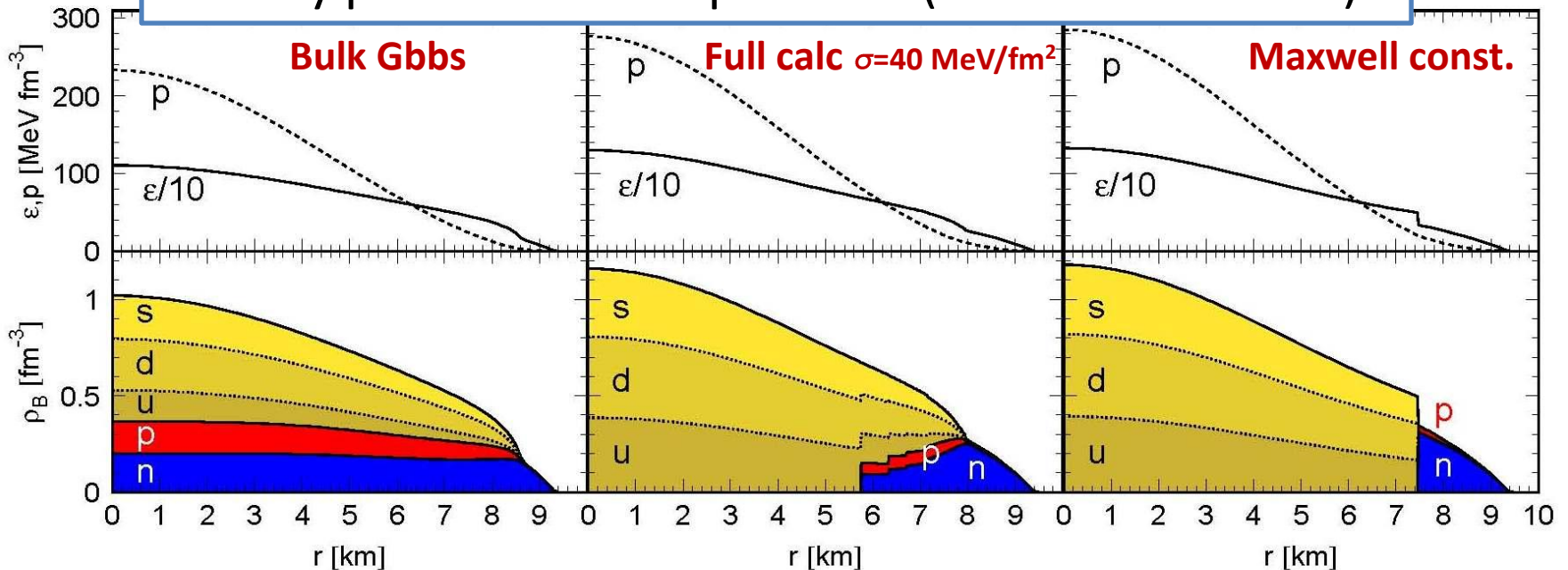
$$m = m(r) = \int_0^r 4\pi s^2 \rho(s) ds$$

Mass inside  $r$

$$M = m(R), \quad R = R(\rho \approx 0)$$

Total mass and Radius

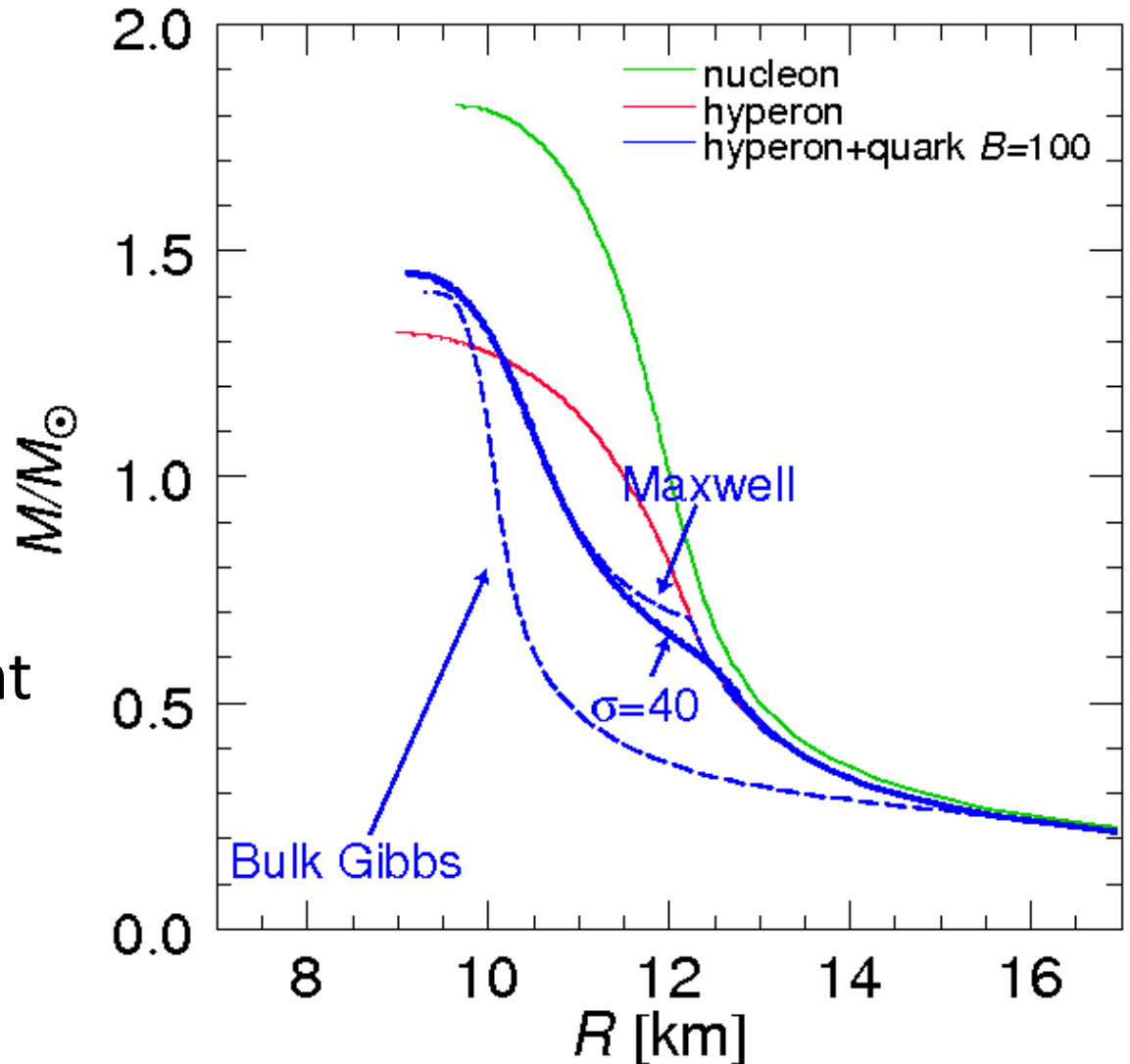
Density profile of a compact star ( $M=1.4$  solar mass)



# Mass-Radius relation of compact stars

Full calc yields the neutron star mass very close to that of the Maxwell constr.

The maximum mass are not very different for three cases.



# Summary

- We have studied “Pasta” structures of Hyperon-quark mixed phase by means of BHF and MIT bag model.
- Resultant EOS of mixed phase is close to that of the Maxwell constrctn instead of a bulk Gibbs calc.
- The mass-radius relation of a compact star is also close to the Maxwell constrctn case.
- But the particle fraction and the inner structure is quite different. Hyperons are strongly suppressed.
  - important for thermal property and  $\nu$  opacity.