

# Role of the explicit tensor correlation in neutron halo nuclei

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1. Tensor correlation in light nuclei
2. LS splitting in He isotopes (poster, QM-050)
3. Disappearance of shell gap at  $N=8$  in Li isotopes

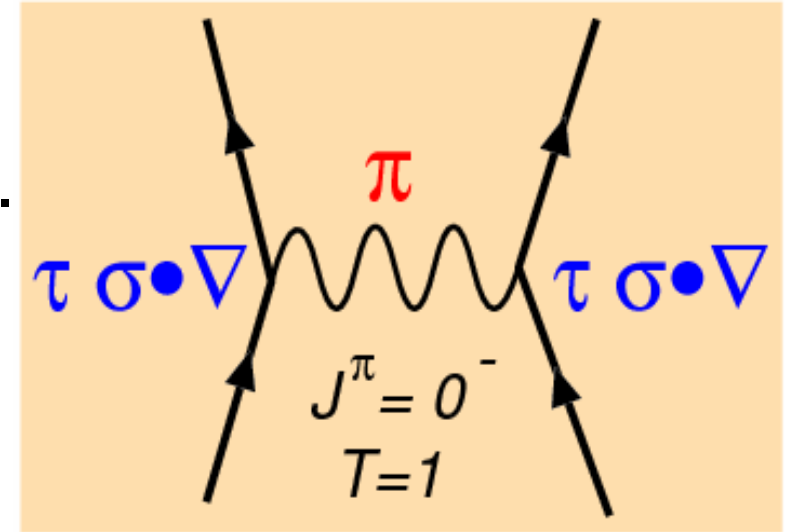
In collaboration with

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Hiroshi Toki (RCNP)    Kiyomi Ikeda (RIKEN)

# Motivation

- Tensor force ( $V_{tensor}$ ) plays a significant role in the nuclear structure.
  - In  ${}^4\text{He}$ ,  $\langle V_{tensor} \rangle \sim \langle V_{central} \rangle$
  - $\frac{V_{\pi}}{V_{NN}} \sim 80\%$  (GFMC)

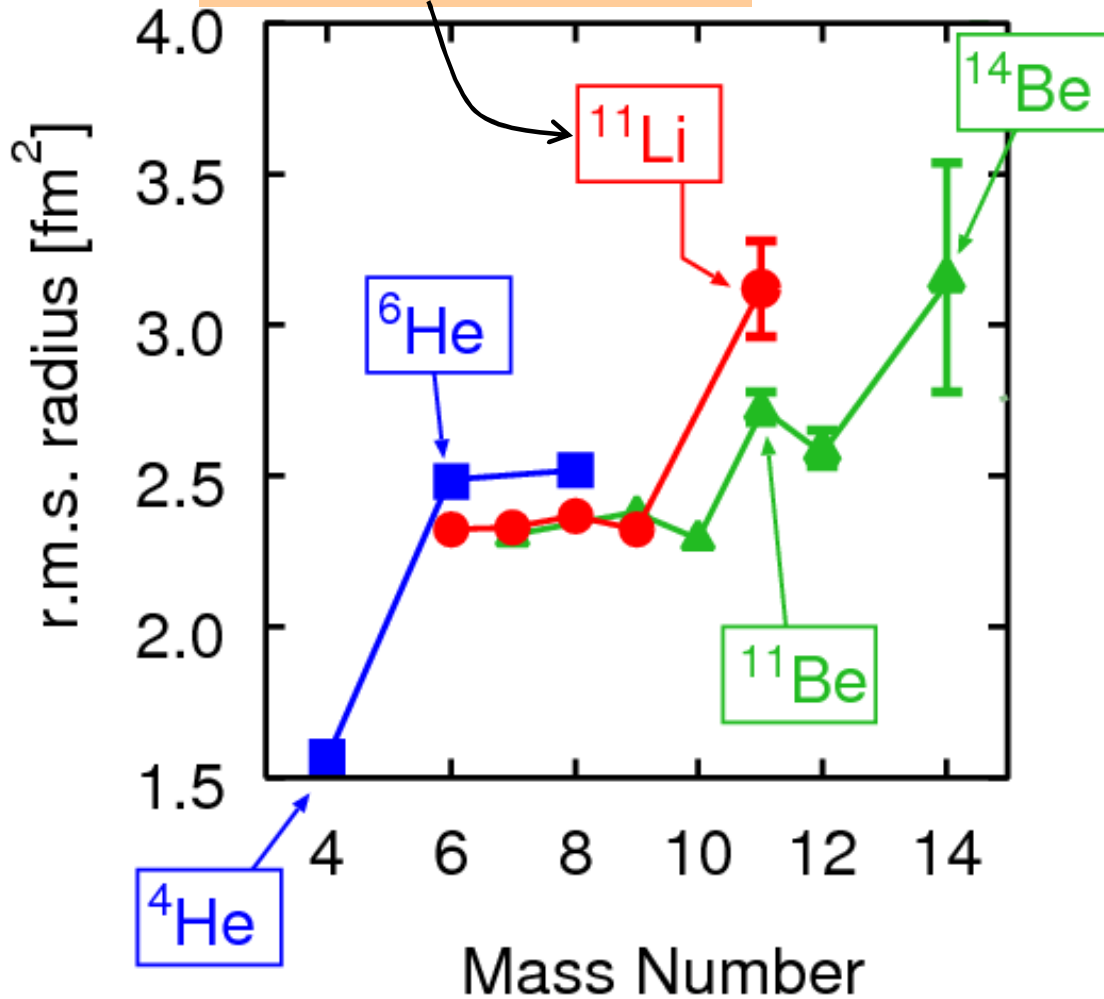


R.B. Wiringa, S.C. Pieper, J. Carlson, V.R. Pandharipande, PRC62(2001)

- We would like to understand the role of  $V_{tensor}$  in the nuclear structure **by describing tensor correlation explicitly.**
  - model wave function (shell model and cluster model).
  - spatial properties, p-h correlation, ...
- Spectroscopy of neutron-rich nuclei : He and Li isotopes

# Characteristics of Li-isotopes

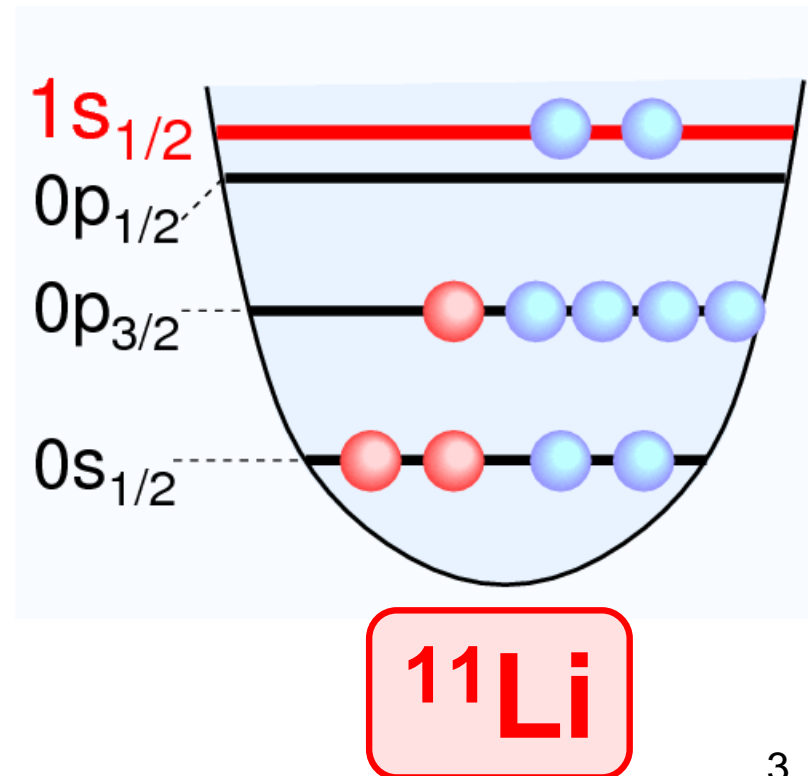
## Halo structure



I. Tanihata et. al  
PLB206(1988)592

## ► Breaking of magicity N=8

- <sup>10-11</sup>Li, <sup>11-12</sup>Be
- <sup>11</sup>Li ... (1s)<sup>2</sup> ~ 50%.  
(Expt by Simon et al., PRL83)
- **Mechanism is unclear**



# $^{11}\text{Li}$ with coupled $^9\text{Li}+n+n$ model

- System is solved based on RGM

$$H(^{11}\text{Li}) = H(^9\text{Li}) + H_{nn} \quad \Phi(^{11}\text{Li}) = \mathcal{A} \left\{ \sum_{i=1}^N \psi_i(^9\text{Li}) \cdot \chi_i(nn) \right\}$$

$$\sum_{i=1}^N \left\langle \psi_j(^9\text{Li}) \left| H(^{11}\text{Li}) - E \right| \mathcal{A} \left\{ \psi_i(^9\text{Li}) \cdot \chi_i(nn) \right\} \right\rangle = 0$$

$\psi_i(^9\text{Li})$ : **shell model type configuration**

- Orthogonality Condition Model (OCM) is applied.

$$\sum_{i=1}^N \left[ h_{ij} (^9\text{Li}) + (T_1 + T_2 + V_{c1} + V_{c2} + V_{12}) \cdot \delta_{ij} \right] \chi_j(nn) = E \chi_i(nn)$$

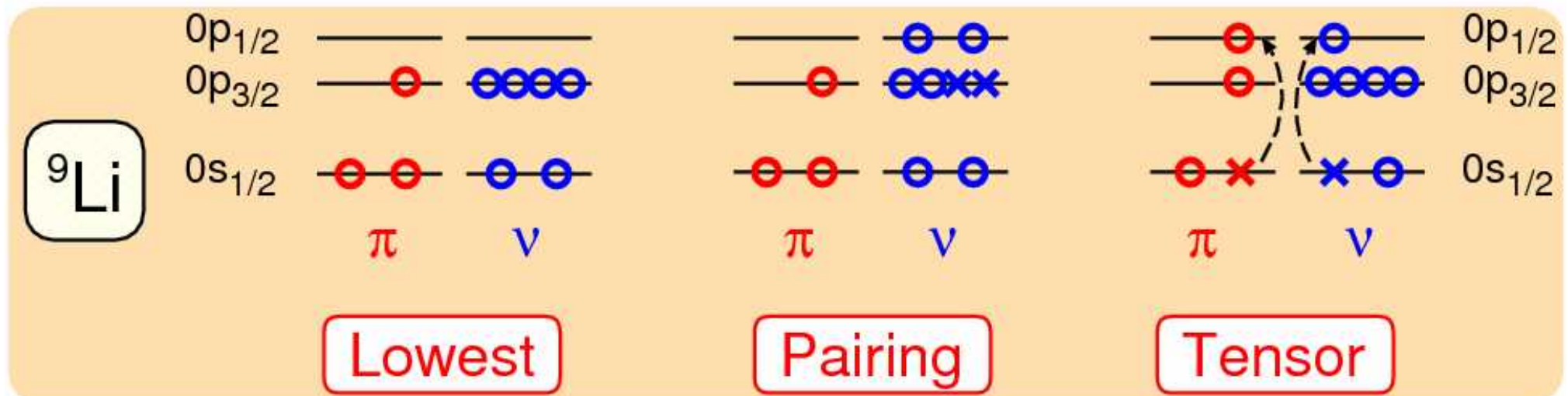
$h_{ij} (^9\text{Li}) = \langle \psi_i | H(^9\text{Li}) | \psi_j \rangle$  : **Hamiltonian for  $^9\text{Li}$**

$\chi(nn) = \mathcal{A} \{ \phi_1 \phi_2 \}$  : **2 neutrons with Gaussian expansion method**

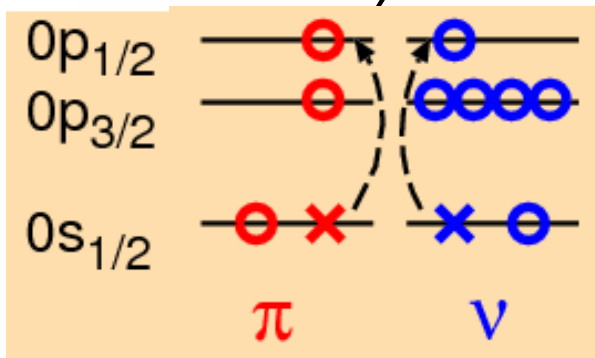
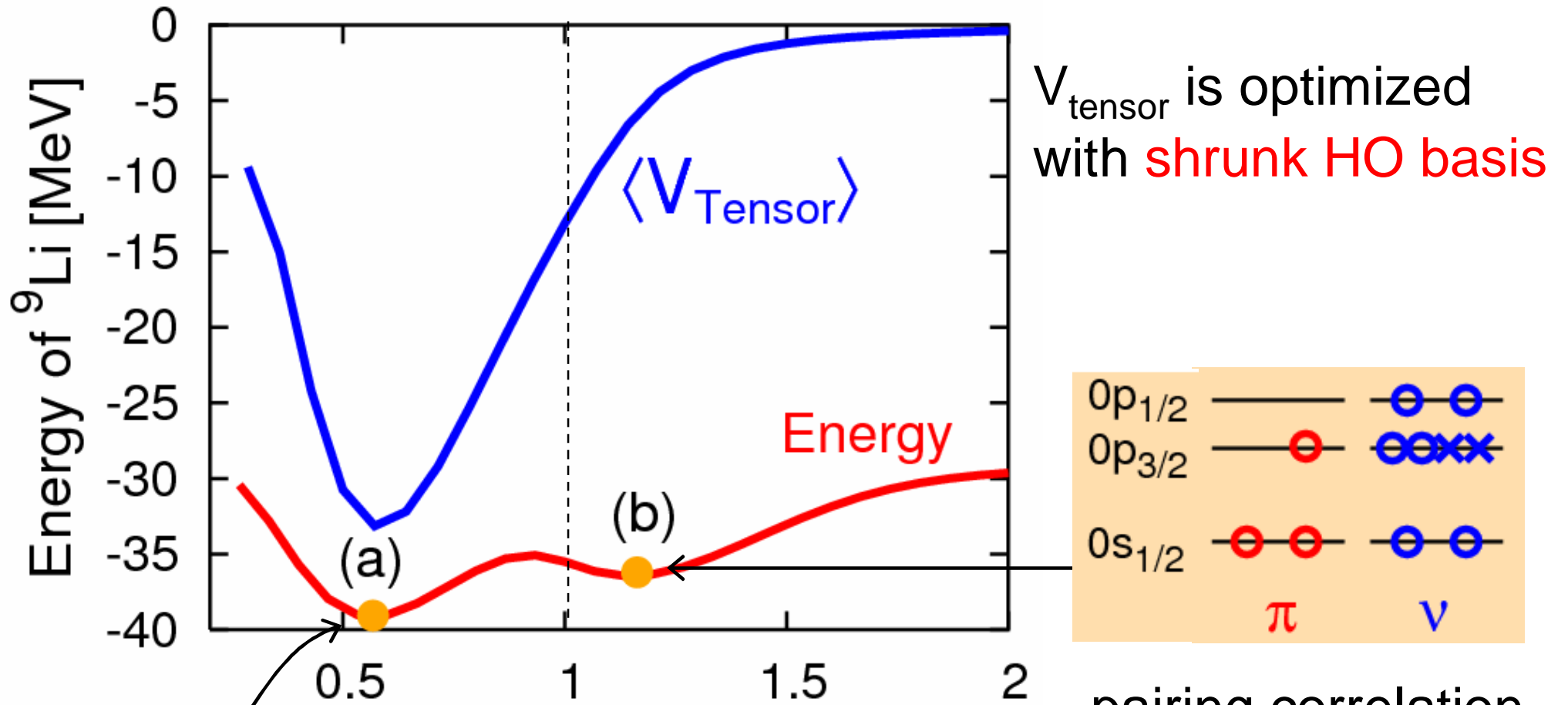
$\langle \phi_i | \phi_\alpha \rangle = 0, \{ \phi_\alpha \in ^9\text{Li} \}$  : **Orthogonality to the Pauli-forbidden states** <sup>4</sup>

# ${}^9\text{Li}$ with tensor and pairing correlations

- Tensor-optimized shell model  
TM et al., PTP113(2005), PTP117(2007), Terasawa PTP22(1969).
- $0s+0p+1s0d$  within  $2p2h$  excitations.
- Length parameters  $b_{0s}$ ,  $b_{0p1/2}$ ,  $b_{0p3/2}$ , ... are determined **independently** and **variationally**.
  - Describe **high momentum component** from  $V_{\text{tensor}}$   
cf. CPP-HF by Sugimoto et al., NPA740 / Akaishi NPA738



# Energy surface for b-parameter in ${}^9\text{Li}$

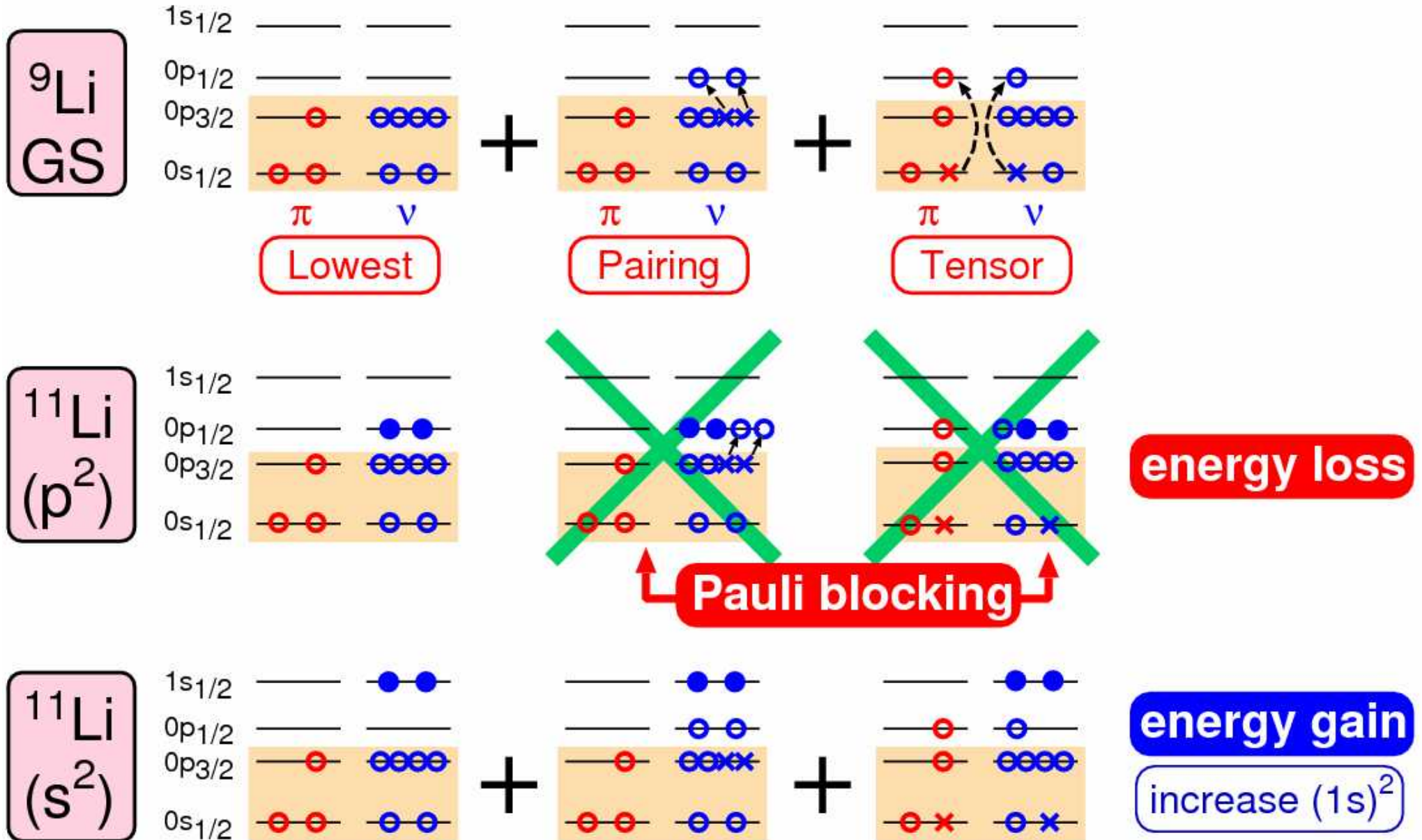


$b_{0p_{1/2}} / b_{0s}$  [fm]

Tensor correlation from the bare  $V_{\text{tensor}}$

pairing correlation of P-shell

# Expected effects of pairing and tensor correlations in $^{11}\text{Li}$

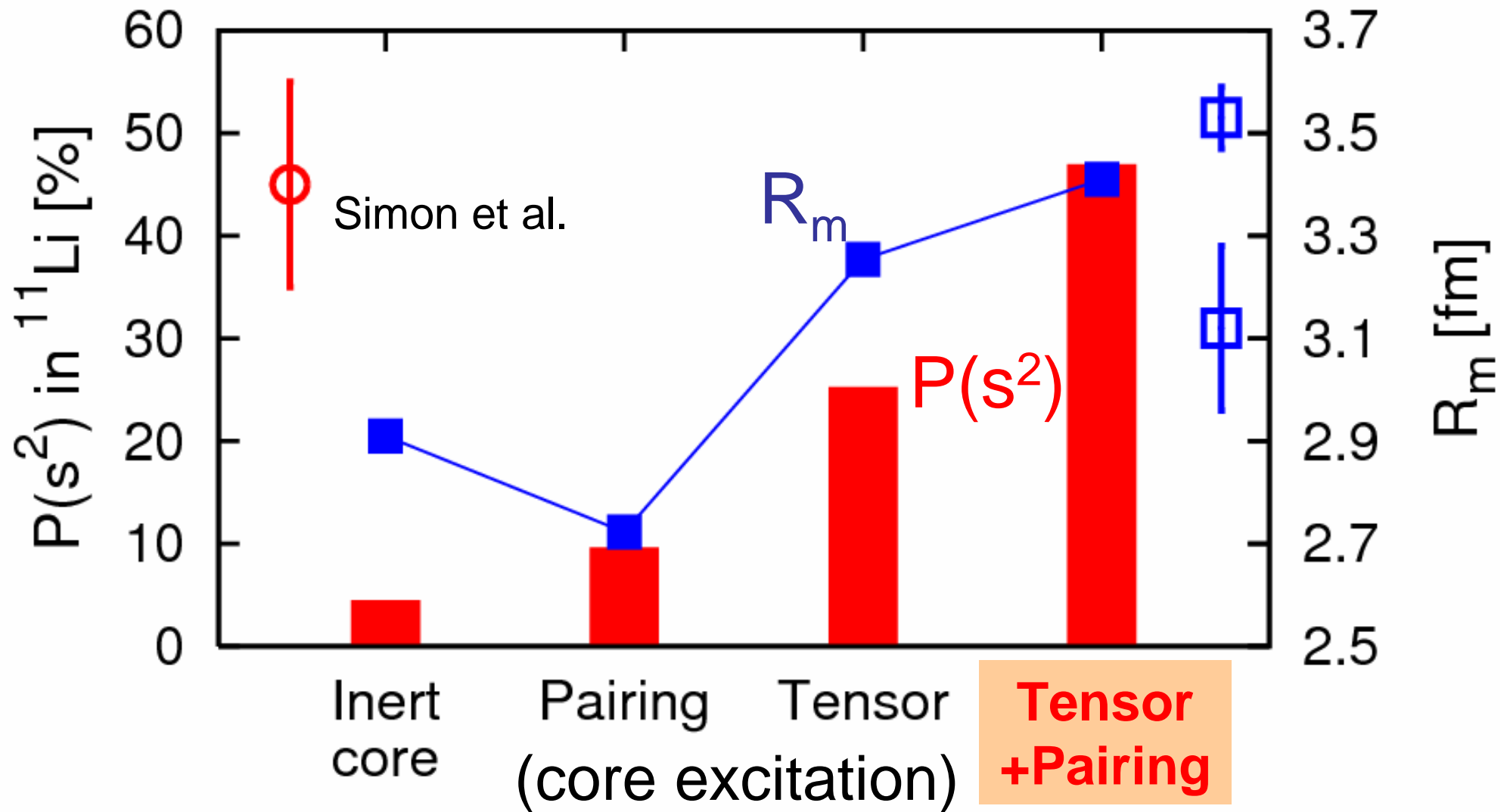


Pairing-blocking :

K.Kato,T.Yamada,K.Ikeda,PTP101('99)119, Masui,S.Aoyama,TM,K.Kato,K.Ikeda,NPA673('00)207.

TM,S.Aoyama,K.Kato,K.Ikeda,PTP108('02)133, H.Sagawa,B.A.Brown,H.Esbensen,PLB309('93)1.

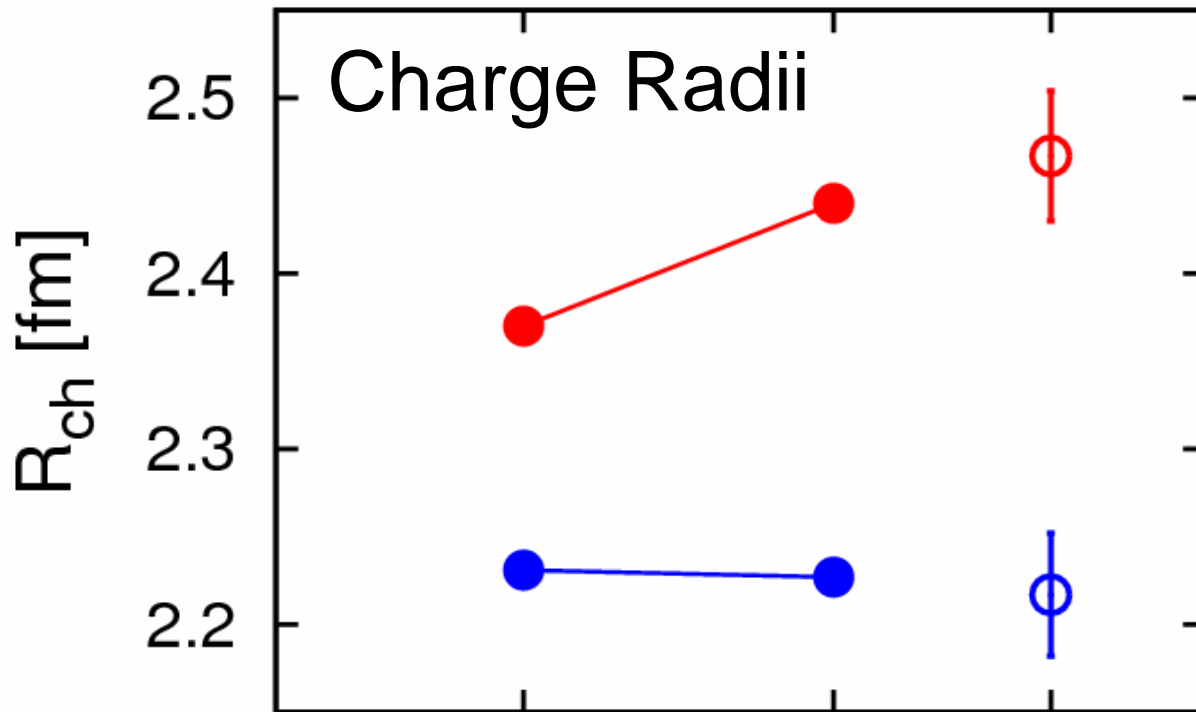
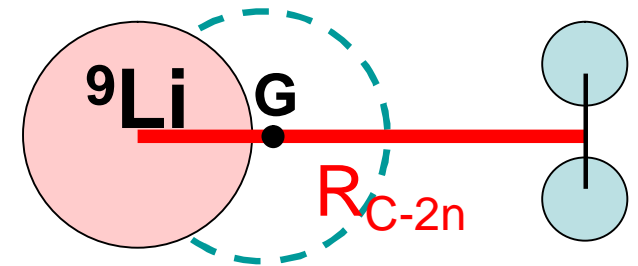
# $^{11}\text{Li}$ G.S. properties ( $S_{2n}=0.31$ MeV)





# Charge Radii of $^{11}\text{Li}$

$$R_{\text{proton}}^2(^{11}\text{Li}) = R_{\text{proton}}^2(^9\text{Li}) + \left(\frac{2}{11}\right)^2 R_{\text{C-2n}}^2$$



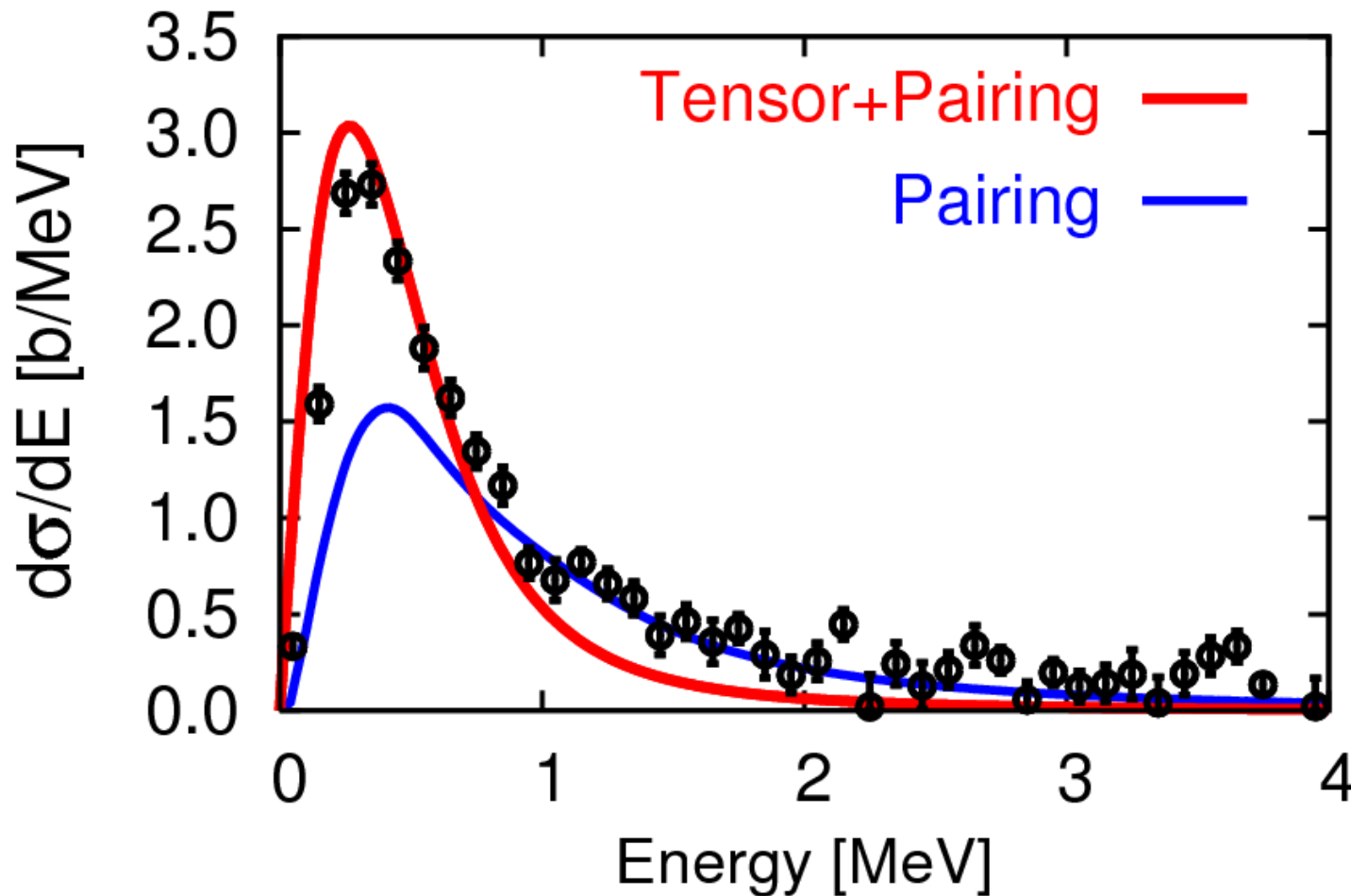
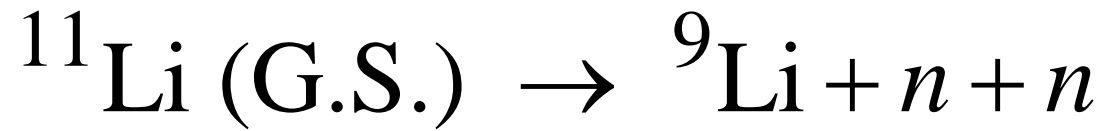
Inert  
core

**Tensor  
+Pairing**

Expt. (Sanchez et al., PRL96(2006))

$R_{\text{C-2n}}$	4.67	<b>5.69</b>	[fm]
$P(s^2)$	4	47	%

# Coulomb breakup strength of $^{11}\text{Li}$



**No three-body  
resonance**

E1 strength by using the  
Green's function method  
+Complex scaling method  
+Equivalent photon method  
(TM et al., PRC63('01))

- Expt: T. Nakamura et al. , PRL96,252502(2006)
- Energy resolution with  $\sqrt{E} = 0.17$  MeV.

# Virtual s-states in $^{10}\text{Li}$

- $1s_{1/2}$  virtual state:

$$(0p_{3/2})_{\pi} (1s_{1/2})_{\nu} \rightarrow 1^{-}, 2^{-}$$

$a_s$  : scattering length of  $^9\text{Li}+n$

**Expt.** M. Thoennesen et al.,  
PRC59 (1999)111.  
M. Chartier et al.  
PLB510(2001)24.  
H.B. Jeppesen et al.  
PLB642(2006)449.

$$a_s = -10 \sim -25 \text{ fm}$$

cf.  $a_s(nn) : -18.5 \pm 0.5 \text{ fm}$

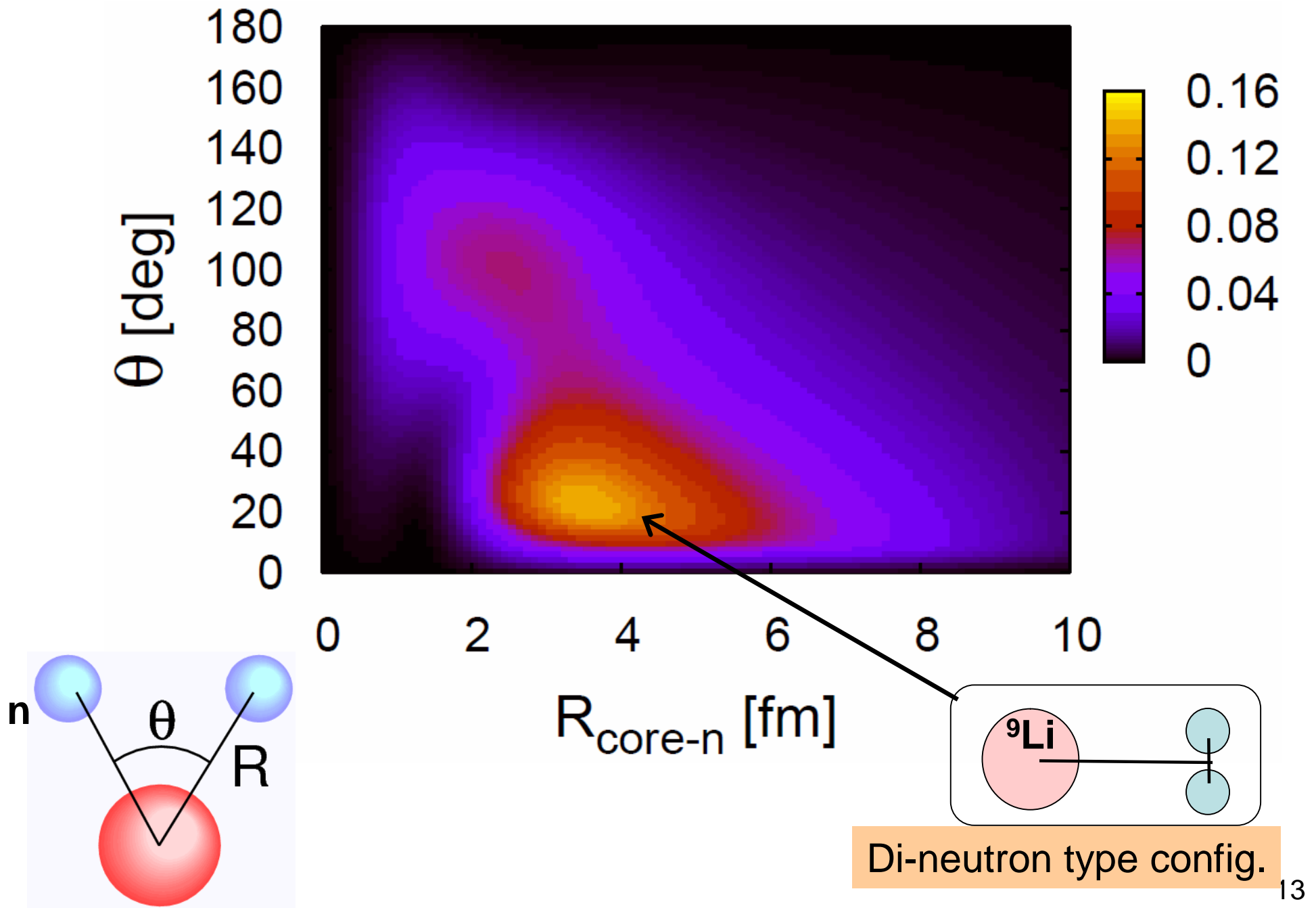
$J^{\pi}$	Inert core	<b>Tensor +Pairing</b>
1 <sup>-</sup>	+1.4 fm	-5.6 fm
2 <sup>-</sup>	+0.8 fm	<b>-17.4 fm</b>

**Pauli-blocking naturally describes virtual s-state in  $^{10}\text{Li}$**

# Summary

- We investigate the explicit tensor correlation in Li-isotopes.
- $^9\text{Li}$  with tensor optimization
  - $0^-$  coupling between  $0s_{1/2}$  and  $0p_{1/2}$ .
  - Deuteron correlation ( $J=1, T=0$ ) for  $2p2h$  excitations.
- $^{10}\text{Li}$ 
  - Pauli-blocking naturally explains the virtual s-states.
- $^{11}\text{Li}$ 
  - Pauli-blocking naturally explains breaking of magicity.
  - Charge radii, Coulomb breakup strength, Q-moment.

# 2n correlation density in $^{11}\text{Li}$



# Pion exchange vs. $V_{\text{Tensor}}$

$$3(\vec{\sigma}_1 \cdot \hat{q})(\vec{\sigma}_2 \cdot \hat{q}) \frac{q^2}{m^2 + q^2} = (\vec{\sigma}_1 \cdot \vec{\sigma}_2) \frac{q^2}{m^2 + q^2} + S_{12} \frac{q^2}{m^2 + q^2}$$

$$= (\vec{\sigma}_1 \cdot \vec{\sigma}_2) \left[ \frac{m^2 + q^2}{m^2 + q^2} - \frac{m^2}{m^2 + q^2} \right] + S_{12} \frac{q^2}{m^2 + q^2}$$

↑  
Delta interaction

Involve large momentum

Tensor operator

$$S_{12} = 3(\vec{\sigma}_1 \cdot \hat{q})(\vec{\sigma}_2 \cdot \hat{q}) - (\vec{\sigma}_1 \cdot \vec{\sigma}_2)$$

↑  
Yukawa interaction

-  $V_{\text{tensor}}$  produces the high momentum component.